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Adsorptive graphene doping: effect of a polymer contaminant

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ABSTRACT:

Transfer-induced contamination of graphene and the limited stability of adsorptive dopants are two of the main issues faced in the practical realization of graphene-based electronics. Herein, we assess the stability of HNO₃, MoO₃, and AuCl₃ dopants upon transferred graphene with different extents of polymer contamination. Sheet resistivity measurements prove polymer residues induce a significant degenerative effect in terms of doping stability for HNO₃ and MoO₃ and a stabilizing effect for AuCl₃. Further characterization by Raman spectroscopy and atomic force microscopy provides insight into the stability mechanism. Together, these findings demonstrate the relevance of contamination in the field of adsorptive doping for the realization of graphene-based functional devices.

Owing to its unique mechanical, electrical, and optical properties, graphene holds much promise for application in electronics, displays, photovoltaics, and energy storage.^[1-4] For example, graphene has been proposed as a cheaper, flexible, and more sustainable alternative to indium tin oxide (ITO) for use in solar cells or touchscreen displays.^[5] One such property is the very high carrier mobility afforded by graphene's unique band structure. Despite this, however, the intrinsic carrier density of pristine graphene is very low, hence doping is required to achieve conductivity values competitive with those of ITO thin films. This has been attempted via substitutional^[6] and adsorptive^[7,8] doping, as well as surface functionalization.^[9] Of these, adsorptive doping shows the most promise. The weak nature of adsorbate binding does not perturb the graphene lattice structure, thus maintaining its carrier mobility. Nevertheless, doping has been found to be unstable both over time and upon heating, limiting the practical application of graphene devices with operating temperatures in the range of 50 - 120 °C.^[10] Given the practical necessity for stable doping, we assess how polymer contamination affects the strength and stability of HNO₃, ^[11] MoO₃,^[12] and AuCl₃^[7] dopants. We find both advantageous and negative effects in terms of dopant stability, and thereby contribute to the understanding of the impact of polymer contaminant on the behavior of adsorptive dopants.

Absorptive polymer impurities commonly result from transfer of graphene during device fabrication, following synthesis via chemical vapor deposition (CVD).^[13] The most popular method involves a wet transfer technique by which a polymer, typically poly(methyl-methacrylate) (PMMA), is coated onto the graphene surface as a supporting layer throughout etchant removal of the deposition catalyst and subsequent transfer to the device substrate.^[14] Dissolution of the PMMA layer then yields, ideally, a pristine graphene surface, but a methodology that facilitates complete PMMA removal without graphene degradation has yet to be realised.^[15,16] This is problematic as remnant polymer residues affect graphene's electronic^[17] and surface^[18] properties.

Herein, we thoroughly assess the effect of graphene contamination on dopant stability via sheet resistivity measurements and atomic force microscopy (AFM). We find that PMMA residues reduce the doping strength and facilitate degradation when HNO₃ and MoO₃ are employed, but induce a beneficial effect for AuCl₃. Following vacuum annealing at elevated temperatures, AuCl₃-doped graphene is up to five times more stable than the HNO₃ and MoO₃ counterparts. On this basis, we observe the need to account for polymer contamination when reporting the performance of adsorptive dopants. This could have wider relevance to other fields where adsorptive doping is applied, such as in the doping of carbon nanotubes or other two dimensional materials.^[19–21]

To obtain these results, we fabricate and measure devices using CVD-grown monolayer graphene, as previously described.^[22] Au/Cr (60:6 nm) electrodes in four-probe geometry, with inter-electrode distances of 2/5/2 µm are deposited on a Si/SiO₂ support (300 nm SiO₂ layer). The graphene is then transferred to Si/SiO₂ substrate through wet transfer with (NH₄)₂S₂O₈ etchant and a supporting PMMA thin film. Graphene

strips 2 µm wide are defined with an O₂ plasma etch through a UVIII resist mask patterned by electron-beam lithography. Prior to doping, adventitious surface adsorbates are removed via an initial vacuum anneal. The sheet resistivity is used as a baseline for subsequent measurements. Doping is achieved by dipping the devices into HNO₃ (68 % aq.) at 50 °C, by spin coating with 20 mM AuCl₃ in MeNO₂ solution, or by thermal evaporation of a 6 nm MoO₃ film at ~10⁻⁶ mbar. The stability of the dopants is assessed by stepwise annealing at 70 °C, 130 °C, and 200 °C at pressures below 3×10^{-6} mbar, with immediate resistivity characterization after each anneal. Currents and voltages are applied/received by source/measurement units under ambient conditions. The maximum applied voltage is 0.06 V and currents are in the range of µA. Sheet resistivity averages and errors are calculated from the measurement of \geq 7 working devices. To expel the PMMA contamination residues from graphene surface,^[15] we employ two regimes: the devices are pre-annealed at < 3 × 10⁻⁶ mbar at either 200 °C for 1 h or 300 °C for 2 hs. The mild pre-annealing conditions partially remove the polymer from the graphene surface giving rise to a thickness (graphene plus contaminants) of 4.3±0.2 nm. In contrast, the stronger conditions give rise to more thorough cleaning, resulting in graphene devices with a thickness of 0.63±0.3 nm.

Following HNO₃ or MoO₃ treatment, we find that the extent of polymer contamination induces a negative effect on the doping strength (Figs 1(a) to 1(d)). Primarily, the absolute sheet resistivity for heavilycontaminated graphene appears to be consistently higher than that for the sparsely-contaminated counterpart (Fig 1(a,c)), regardless of the fact that PMMA mildly p-dopes graphene.^[23] Secondly, the doping of heavilycontaminated graphene is significantly less stable than that of sparsely-contaminated graphene (Fig. 1(b, d)). This is quantified by plotting the average ratio of sheet resistance to as-doped sheet resistance, measured between the same devices at different annealing steps. We evidence 62% and 50% disparities in resistivity increase following annealing at 70 °C, between heavily-contaminated and sparsely-contaminated graphene for MoO₃ and HNO₃, respectively. Furthermore, the instability of heavily-contaminated relative to sparselycontaminated graphene is more pronounced for MoO₃ than for HNO₃ dopant. Over the annealing stages, we observe average differences of 53% and 37% in resistivity between heavily-contaminated and sparselycontaminated graphene for MOO_3 and HNO_3 , respectively. Finally, the reduced homogeneity in dopant stabilities is also evidenced by larger standard deviations in sheet resistivity ratios to as-doped graphene for heavily-contaminated graphene, with a more pronounced effect for MoO₃ dopant. In this case, the average standard deviations across the annealing steps were ± 0.22 and $\pm 0.09 \ \Omega/\Box$ for heavily- and sparselycontaminated graphene, respectively.

The effect of polymer contamination appears to be dramatically different with AuCl₃ dopant (Fig. 1 (e, f)). Following doping, we observe a remarkably low average sheet resistance of $102 \pm 4.5 \ \Omega/\Box$ (for sparsely-contaminated graphene), surpassing that observed in similar studies (Fig. 1(e)).^[24–26] The lowest sheet resistivity observed in a *single* device is 85.8 Ω/\Box . In a clear contrast to HNO₃ and MoO₃ dopants, the stability of the AuCl₃ doping for both heavily-contaminated and sparsely-contaminated graphene remains practically unchanged after annealing at 70 °C and 130 °C. After annealing at 200 °C, however, the heavily-contaminated

graphene samples exhibit greater dopant stability than that on sparsely-contaminated graphene (Fig 1.(f)).

These results are also reflected by Raman characterization (514 nm laser at ~500 μ W with a 50× objective and spot diameter of ~1.5 μ m) of heavily-contaminated graphene before and after doping with all dopants and subsequent annealing at 200 °C (Fig. 2). The un-doped graphene spectrum has an I_{2D}/I_G ratio of ~2.3, a G peak full-width at half maximum of ~30 cm⁻¹ and a negligible D peak, concurrent with high quality, monolayer graphene.^[27] The presence of an additional shoulder at slightly lower wavenumber than the G peak is assigned to carbon contamination of the graphene layer. The fact this is absent in the spectra for HNO₃ and AuCl₃ dopants suggests that the doping process removes some contaminants from the graphene, whilst its presence in the spectra for MoO₃ dopant helps explain the highly unstable MoO₃ doping observed. Additionally, the absence of a notable D peak in the doped spectra indicates that the doping and annealing processes do not induce defects on the graphene sheets.^[25,28,29] A shift to higher wavenumber in the G and 2D peaks is observed for the doped graphene spectra, alongside a variation in relative 2D/G intensities to I_{2D}/I_G ~1. These changes are indicative of the doping effect induced by all three of the dopants investigated.^[30]

To rationalize these observations, we conducted further analysis via AFM (Figs. 3(a)-(f)). We find that after HNO₃ doping and annealing, the heavily-contaminated graphene samples are decorated with large features which we attribute to polymer agglomeration (Fig. 3(c)). In contrast, sparsely-contaminated graphene appeared nearly unchanged from the pre-doped form (Fig. 3(d)). These results evidence that large quantities of PMMA in close contact to the graphene surface display the greatest resilience to displacement by chemical or thermal means, as suggested previously.^[15] We compare these findings to those of previous studies in our laboratory and elsewhere^[28,29,31] concerning MoO₃ and HNO₃ dopants on *pristine* graphene, which report markedly smaller doping instability. In those accounts, MoO₃ was found to be completely stable,^[29] and HNO₃ exhibited only a 2% increase in resistivity following annealing at 130 °C.^[28] Therefore, we resolve that polymer contamination has a significant, negative effect on the stability of MoO₃ and HNO₃ dopants.

For the AuCl₃ case (Fig. 3 (e, f)), we find the samples are decorated with Au nanoparticles (AuNP) which have formed upon reduction of Au(III) species on the graphene surface and melted and agglomerated into larger clusters upon annealing.^[26,32] The number density of nanoparticles is $5.3 \pm 0.5 \ \mu m^{-2}$ for heavily-contaminated graphene and $3.2 \pm 0.2 \ \mu m^{-2}$ for the sparsely-contaminated samples. The particle size (measured above a 15 nm height threshold to account for particles of polymer residue) averages 63 ± 2 and 80 ± 3 nm, respectively. On this basis, we hypothesize that on average, smaller AuNPs are formed on heavily-contaminated graphene as the residues limit Au surface diffusion, in agreement with previous observations.^[33-36] We note that in addition to purely metallic Au on the graphene surface, AuCl₄⁻ and Cl-containing species are also present in the doping process.^[26] Charge-transfer equilibrium between electron-deficient graphene and Cl⁻ produces Cl species, which adsorb to the surface and play a key role in maintaining the doping effect: temperature-induced desorption of Cl₂ results in increased sheet resistivity.^[25,26] In light of the differential AuNP agglomeration, we attribute the stabilizing behavior of the polymer contaminant to reduced surface diffusion of dopant species across the graphene. Such stabilization retards the formation and desorption of Cl₂, and thus slows any increases in resistivity during annealing. Figure 4 cartoons the possible scenarios for

absorptive doping on polymer-contaminated graphene.

This work exemplifies the extent to which polymer residues present a challenge to the control of adsorptively-doped polymer-transferred CVD graphene. Furthermore, the mechanism by which contamination affects dopant stability is not generic and depends upon the nature of the dopant employed. However, we would expect to observe similar, destabilising, behaviour as in the MoO₃ case for other vapour-deposited, adlayer dopants such as WO₃, due to analogous degradation of the dopant/graphene interface. Whilst dopant behaviour is difficult to predict, we believe that polymer contamination can be expected to modify the properties of other p-type dopants such as I₂, and indeed adsorptive n-type dopants also.^[37,38]

In summary, we have explored the effect of a polymer contaminant on the strength and stability of adsorptive graphene dopants. The quantity of contaminating PMMA residue on graphene has been found to negatively affect the stability of HNO₃ and MoO₃-doped graphene at elevated temperatures, with this response most pronounced for MoO₃. Conversely, when doped with AuCl₃, a stabilising effect is recorded at high temperatures in the presence of heavier polymer contamination. These findings are significant to the production and accurate characterisation of stable, adsorptively-doped graphene.

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- [1] M. F. El-Kady, R. B. Kaner, ACS Nano 2014, 8, 8725–9.
- [2] S. Bae, H. Kim, Y. Lee, X. Xu, J.-S. Park, Y. Zheng, J. Balakrishnan, T. Lei, H. R. Kim, Y. Il Song, et al., *Nat. Nanotechnol.* **2010**, *5*, 574–8.
- [3] C.-T. Chien, P. Hiralal, D.-Y. Wang, I.-S. Huang, C.-C. Chen, C.-W. Chen, G. A. J. Amaratunga, *Small* **2015**, *11*, 2929–37.
- [4] L. Kavan, J.-H. Yum, M. Grätzel, Nano Lett. 2011, 11, 5501–6.
- [5] G. Jo, M. Choe, S. Lee, W. Park, Y. H. Kahng, T. Lee, *Nanotechnology* **2012**, *23*, 112001.
- [6] D. Wei, Y. Liu, Y. Wang, H. Zhang, L. Huang, G. Yu, *Nano Lett.* **2009**, *9*, 1752–8.
- [7] K. C. Kwon, B. J. Kim, J.-L. Lee, S. Y. Kim, J. Mater. Chem. C 2013, 1, 2463–2469.
- [8] T. O. Wehling, K. S. Novoselov, S. V Morozov, E. E. Vdovin, M. I. Katsnelson, A. K. Geim, A. I. Lichtenstein, *Nano Lett.* 2008, 8, 173–7.
- [9] A. Sinitskii, A. Dimiev, D. A. Corley, A. A. Fursina, D. V Kosynkin, J. M. Tour, ACS Nano 2010, 4, 1949– 54.
- [10] H. A. Koffi, V. C. K. Kakane, A. Kuditcher, A. F. Hughes, M. B. Adeleye, J. K. . Amuzu, *African J. Sci. Technol. Innov. Dev.* 2015, 7, 485–490.
- [11] H.-J. Lee, E. Kim, J. Park, W. Song, K.-S. An, Y. S. Kim, J.-G. Yook, J. Jung, Carbon N. Y. 2014, 78, 532– 539.
- [12] P. R. Kidambi, C. Weijtens, J. Robertson, S. Hofmann, J. Meyer, Appl. Phys. Lett. 2015, 106, 63304.
- [13] R. Muñoz, C. Gómez-Aleixandre, Chem. Vap. Depos. 2013, 19, 297–322.
- [14] X. Li, W. Cai, J. An, S. Kim, J. Nah, D. Yang, R. Piner, A. Velamakanni, I. Jung, E. Tutuc, et al., Science 2009, 324, 1312–4.
- [15] Y. C. Lin, C. C. Lu, C. H. Yeh, C. Jin, K. Suenaga, P. W. Chiu, Nano Lett. 2012, 12, 414–419.
- [16] X. Liang, B. A. Sperling, I. Calizo, G. Cheng, C. A. Hacker, Q. Zhang, Y. Obeng, K. Yan, H. Peng, Q. Li, et al., ACS Nano 2011, 5, 9144–9153.
- [17] A. Pirkle, J. Chan, A. Venugopal, D. Hinojos, C. W. Magnuson, S. McDonnell, L. Colombo, E. M. Vogel, R. S. Ruoff, R. M. Wallace, *Appl. Phys. Lett.* 2011, *99*, 122108.
- [18] M. Kratzer, B. C. Bayer, P. R. Kidambi, A. Matković, R. Gajić, A. Cabrero-Vilatela, R. S. Weatherup, S. Hofmann, C. Teichert, *Appl. Phys. Lett.* **2015**, *106*.
- [19] R. Graupner, J. Abraham, A. Vencelová, T. Seyller, F. Hennrich, M. M. Kappes, A. Hirsch, L. Ley, *Phys. Chem. Chem. Phys.* **2003**, *5*, 5472–5476.
- [20] K. K. Kim, J. J. Bae, H. K. Park, S. M. Kim, H.-Z. Geng, K. A. Park, H.-J. Shin, S.-M. Yoon, A. Benayad, J.-Y. Choi, et al., J. Am. Chem. Soc. 2008, 130, 12757–12761.
- [21] S. Esconjauregui, L. D'Arsié, Y. Guo, J. Yang, H. Sugime, S. Caneva, C. Cepek, J. Robertson, ACS Nano 2015, 9, 10422–10430.
- [22] X. Wu, G. Zhong, L. D'Arsié, H. Sugime, S. Esconjauregui, A. W. Robertson, J. Robertson, *Sci. Rep.* 2016, 6, 21152.
- [23] A. Pirkle, J. Chan, A. Venugopal, D. Hinojos, C. W. Magnuson, S. McDonnell, L. Colombo, E. M. Vogel, R. S. Ruoff, R. M. Wallace, *Appl. Phys. Lett.* 2011, *99*, 122108.
- [24] K. C. Kwon, B. J. Kim, J.-L. Lee, S. Y. Kim, J. Mater. Chem. C 2013, 1, 2463.
- [25] K. W. Kim, W. Song, M. W. Jung, M.-A. Kang, S. Y. Kwon, S. Myung, J. Lim, S. S. Lee, K.-S. An, *Carbon N. Y.* 2015, 82, 96–102.
- [26] D. Hee Shin, J. Min Kim, C. Wook Jang, J. Hwan Kim, S. Kim, S. H. Choi, J. Appl. Phys. 2013, 113, 3–7.
- [27] Y. Ying Wang, Z. Hua Ni, T. Yu, Z. X. Shen, H. Min Wang, Y. Hong Wu, W. Chen, A. T. Shen Wee, J. Phys. Chem. C 2008, 112, 10637–10640.
- [28] L. D'Arsié, S. Esconjauregui, R. S. Weatherup, X. Wu, W. E. Arter, H. Sugime, C. Cepek, J. Robertson, *RSC Adv.* 2016, 6, 113185–113192.
- [29] L. D'Arsié, S. Esconjauregui, R. Weatherup, Y. Guo, S. Bhardwaj, A. Centeno, A. Zurutuza, C. Cepek, J. Robertson, *Appl. Phys. Lett.* 2014, 105, 2012–2017.
- [30] A. Das, S. Pisana, B. Chakraborty, S. Piscanec, S. K. Saha, U. V Waghmare, K. S. Novoselov, H. R. Krishnamurthy, A. K. Geim, A. C. Ferrari, et al., *Nat. Nanotechnol.* 2008, *3*, 210–5.
- [31] Z. Chen, I. Santoso, R. Wang, L. F. Xie, H. Y. Mao, H. Huang, Y. Z. Wang, X. Y. Gao, Z. K. Chen, D. Ma, et al., *Appl. Phys. Lett.* 2010, 96, 213104.
- [32] R. A. Nistor, D. M. Newns, G. J. Martyna, ACS Nano 2011, 5, 3096–3103.
- [33] L. Meli, P. F. Green, ACS Nano 2008, 2, 1305–1312.
- [34] R. Popescu, R. Schneider, D. Gerthsen, A. Böttcher, D. Löffler, P. Weis, M. M. Kappes, *Surf. Sci.* 2009, 603, 3119–3125.
- [35] M. Watanabe, T. Akimoto, E. Kondoh, *Phys. Status Solidi A* 2012, 209, 2514–2520.
- [36] S. B. Simonsen, I. Chorkendorff, S. Dahl, M. Skoglundh, J. Sehested, S. Helveg, J. Catal. 2011, 281, 147– 155.

- [37] B. H. Kim, S. J. Hong, S. J. Baek, H. Y. Jeong, N. Park, M. Lee, S. W. Lee, M. Park, S. W. Chu, H. S. Shin, et al., Sci. Rep. 2012, 2, 666–669.
- [38] I. Jo, Y. Kim, J. Moon, S. Park, J. S. Moon, W. B. Park, J. S. Lee, B. H. Hong, K. S. Novoselov, A. K. Geim, et al., *Phys. Chem. Chem. Phys.* **2015**, *17*, 29492–29495.

FIG. 1: Sheet resistivity of heavily- and sparsely-contaminated graphene following initial anneal, doping, and annealing at 70 °C, 130 °C and 200 °C for (a) HNO_3 , (c) MoO_3 and (e) $AuCl_3$. Ratios of sheet resistance following annealing for heavily- and sparsely-contaminated graphene relative to that of as-doped graphene for (b) HNO_3 , (d) MoO_3 and (f) $AuCl_3$.

FIG. 2: Raman spectrum of as-transferred heavily-contaminated graphene, and spectra following doping by HNO₃, MoO₃ and AuCl₃ and annealing at 200 °C.

FIG. 3: AFM micrographs of (a) heavily and (b) sparsely contaminated graphene before doping, and following doping by HNO₃ and AuCl₃ and annealing at 200 °C, (c, d) and (e, f), respectively. Note differing height scales between doping conditions. Lateral scale bar = 1 μ m.

FIG. 4: Cartoon showing variation in AuNP growth, dependent on the extent of polymer contamination for heavily-contaminated (a) and sparsely-contaminated graphene (b).





FIG. 2









